

OTTICA @

24/11/20

rappres differenziali delle mappe operatoriali:

$$D_s(a \cdot) = \alpha + \frac{1-s}{2} \frac{\partial}{\partial \alpha^*}$$

$$D_s(a^+ \cdot) = \alpha - \frac{1+s}{2} \frac{\partial}{\partial \alpha}$$

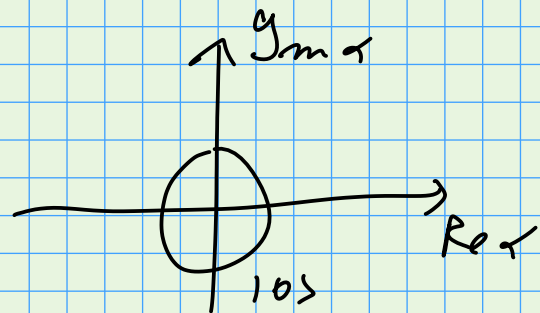
$$D_s(\cdot a) = \alpha - \frac{1+s}{2} \frac{\partial}{\partial \alpha^*}$$

$$D_s(\cdot a^+) = \alpha^* + \frac{1-s}{2} \frac{\partial}{\partial \alpha}$$

Fn di Wigner dello st coerente:

$$W(\alpha, |\beta\rangle\langle\beta|) = \frac{2}{\pi} e^{-2|\alpha - \beta|^2}$$

$$\Rightarrow W(\alpha, |0\rangle\langle 0|) = \frac{2}{\pi} e^{-2|\alpha|^2}$$



$$W(\alpha, |\beta\rangle\langle\beta|) = \int \frac{d^2\lambda}{\pi^2} e^{\alpha\lambda^* - \alpha^*\lambda} \langle \beta | D(\lambda) | \beta \rangle =$$

$$\langle \beta | \beta + \lambda \rangle = e^{-\frac{|\lambda|^2}{2}} e^{\frac{1}{2}(\lambda\beta^* - \lambda^*\beta)}$$

$$= \int \frac{d^2\lambda}{\pi^2} e^{\lambda^*(\alpha - \beta) - \lambda(\alpha^* - \beta^*)} e^{-\frac{|\lambda|^2}{2}} = \frac{2}{\pi} e^{-2|\alpha - \beta|^2}$$

$$\int d^2\lambda e^{\alpha\lambda^* - \lambda\alpha^*} e^{-\frac{|\lambda|^2}{2}} = e^{-6|\alpha|^2}$$

## STATO di FOCK, $|n\rangle$

è auto st dell'Hamiltoniana di singolo modo

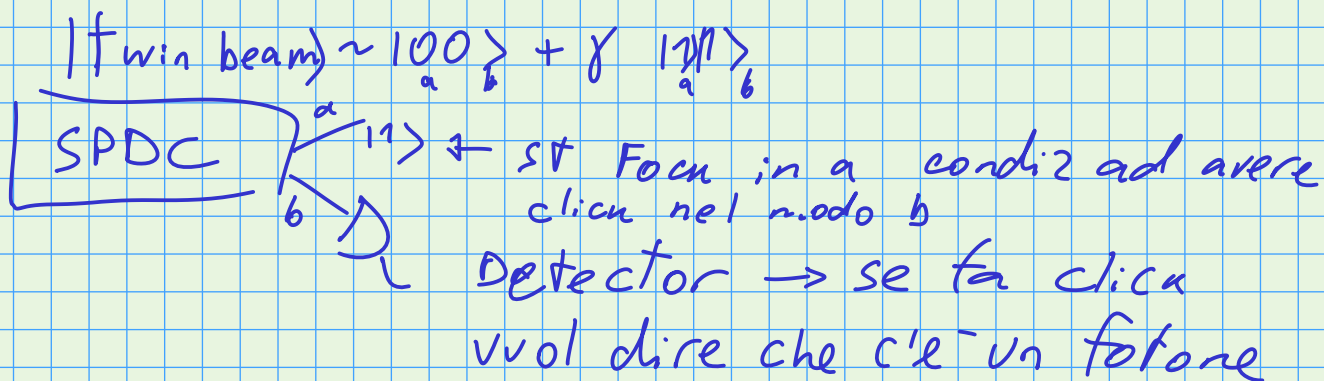
$$a^\dagger a |n\rangle = n |n\rangle$$

fn di Wigner  $\rightarrow$  polinomi di Laguerre

$|0\rangle$  è st coerente e st di Fock

tutti gli altri non sono st a minima indet

Trucco per creare st di Fock  $|1\rangle$

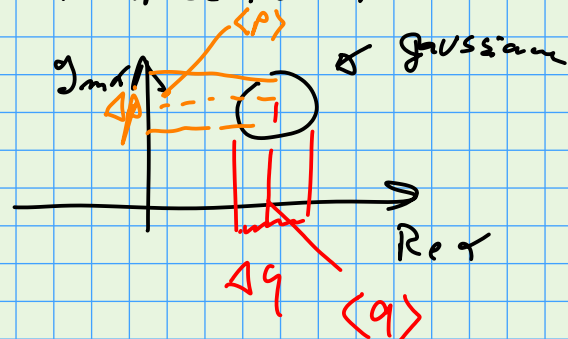


## STATO SQUEEZED

abbiamo visto st coerenti sono a minima indet

$$\Delta p \Delta q = \frac{1}{2} \quad \text{ma} \quad \Delta p = \Delta q = \frac{1}{\sqrt{2}}$$

$\uparrow$  quadrature  $\perp$



lo stato squeezed ha quadrature diverse

st è squeezed  $\stackrel{\text{def}}{\Leftrightarrow} \exists A$  t.c.  $\Delta A^2 < \frac{1}{2} |\langle [A, B] \rangle|$

Heis-Robertson

$$\Delta A^2 \Delta B^2 \geq \frac{1}{4} |\langle [A, B] \rangle|^2$$

st squeezed ideale  $\stackrel{\text{def}}{=} \text{st a minima indet}$

$$\Delta A^2 \Delta B^2 = \frac{1}{4} |\langle [A, B] \rangle|^2$$

op di SQUEEZING  $S(\zeta) \stackrel{\text{def}}{=} e^{\frac{1}{2}(\zeta^* a^2 - \zeta a^{+2})}$

e' op unitario  $S^\dagger = S^{-1}$ : e' un op evoluzione unitaria  $e^{-\frac{i}{\hbar} H t}$  con  $H \propto i(\zeta^* a^2 - \zeta a^{+2})$

Hamilt di Squeezing  $\rightarrow$  cristallo non lineare  $\chi^{(2)}$

$$\begin{array}{c} |a\rangle_c \\ \hline a=b \\ \hline \end{array} \boxed{H = \chi^2 (a^{+2} + a^2)} \xrightarrow{|a\rangle_c} \text{a st squeezed}$$

$S(\zeta) |0\rangle = | \text{vuoto squeezed} \rangle$

BCH di  $SU(1,1)$ :  $K_+ \stackrel{\text{def}}{=} \frac{a^{+2}}{2} = K_-^\dagger$

$$K_z \stackrel{\text{def}}{=} \frac{1}{2} (a^\dagger a + \frac{1}{2})$$

soddisfa algebra di  $SU(1,1)$

$$[K_+, K_-] = \frac{1}{4} [a^{+2}, a^2] = \frac{1}{4} (a^\dagger [a^\dagger, a^2] + [a^\dagger, a^2] a^\dagger) =$$

$$= \frac{1}{4} (a^\dagger a [a^\dagger, a] + a^\dagger [a^\dagger, a] a + [a^\dagger, a] a a^\dagger + a [a^\dagger, a] a^\dagger)$$

$$= -\frac{1}{4} (a^\dagger a + a^\dagger a + a a^\dagger + a a^\dagger) = -a^\dagger a - \frac{1}{2} = -2K_z$$

$$[K_z, K_\pm] = K_\pm$$

$$\Rightarrow \text{BCH di } \text{su}(1,1) \quad S(\xi) = e^{\frac{\xi^+ a^2}{2}} e^{-\frac{\xi}{2} a^{\dagger 2}} \stackrel{\text{BCH}}{=} e^{-\frac{\xi}{2} a^{\dagger 2}} \left( \text{ch} \left| \frac{\xi}{2} \right| \right)^{-a^{\dagger} a - \frac{1}{2}} e^{\frac{\xi^+ a^2}{2}}$$

$$\xi = \frac{\xi}{|\xi|} \text{th} |\xi|$$

$\xi$  def PARAMETRO di SQUEEZING EC

VUOTO SQUEEZED

$$S(\xi) |0\rangle = e^{-\frac{\xi a^{\dagger 2}}{2}} \left( \text{ch} \left| \frac{\xi}{2} \right| \right)^{-a^{\dagger} a - \frac{1}{2}} e^{\frac{\xi^+ a^2}{2}} |0\rangle$$

$$= \frac{1}{\sqrt{\text{ch} |\xi|}} \sum_{n=0}^{\infty} \frac{(-\xi)^n}{n!} \underbrace{a^{\dagger 2n} |0\rangle}_{\sqrt{(2n)!} |2n\rangle} = \frac{1}{\sqrt{\text{ch} |\xi|}} \sum_{n=0}^{\infty} \sqrt{\binom{2n}{n}} (-\xi)^n |2n\rangle$$

$\frac{\sqrt{(2n)!}}{n!} = \sqrt{\frac{(2n)!}{n! n!}}$   
 solo i termini pari nelle base di Fock

il vuoto squeezed ha energia:

$$\langle a^{\dagger} a \rangle = \langle 0 | S^{\dagger} a^{\dagger} a S |0\rangle = \text{sh}^2 |\xi| = 0 \Leftrightarrow \xi = 0$$

squeezed coherent,  $t$  st  $S(\xi) |\alpha\rangle = S(\xi) D(\alpha) |0\rangle$

displaced squeezed vacuum  $D(\alpha) S(\xi) |0\rangle$

VARIANZA delle QUADRATURE del VUOTO SQUEEZED

→ vedremo che vuoto squeezed e  $s^{\dagger}$  squeezed

" " e  $s$  a minima indet

$$\Delta q \Delta p = \frac{1}{2} \quad \Delta q = \frac{e^{-|\xi|}}{\sqrt{2}} < \frac{1}{\sqrt{2}} = \sqrt{\frac{1}{2} |\langle [A, B] \rangle|} = \frac{1}{\sqrt{2}}$$

$$\Delta p = \frac{e^{+|\beta|}}{\sqrt{\Sigma}}$$

$$\Delta a_\varphi^2 = \langle a_\varphi^2 \rangle - \langle a_\varphi \rangle^2$$

$$a_\varphi = \frac{a^\dagger e^{i\varphi} + a e^{-i\varphi}}{2}$$

$$q = a_{\varphi=0} \quad p = a_{\varphi=\frac{\pi}{2}}$$

Schrödinger

Heis

$$\langle X \rangle = (\langle 0 | S^\dagger) X (S | 0 \rangle) = \langle 0 | (S^\dagger X S) | 0 \rangle$$

$$S^\dagger a_\varphi S = \frac{S^\dagger a^\dagger S e^{i\varphi} + S^\dagger a S e^{-i\varphi}}{2}$$

$$e^{\uparrow} B e^{-A} = B + [A, B] + \dots + \frac{1}{n!} [A, [A, \dots [A, B]]] + \dots$$

$$S^\dagger a S = a + \left[ -\frac{\xi^* a^2}{2} + \frac{\xi a^{\dagger 2}}{2}, a \right] + \frac{1}{2!} \left[ -\frac{\xi^* a^2}{2} + \frac{\xi a^{\dagger 2}}{2}, -\xi a^\dagger \right]$$

$$S^\dagger = e^{-\frac{1}{2}(\xi^* a^2 - \xi a^{\dagger 2})}$$

$$-\xi a^\dagger$$

$$[B, [B, A]]$$

$$+ \frac{1}{3!} \left[ -\frac{\xi^* a^2}{2} - \frac{\xi a^{\dagger 2}}{2}, \xi | \xi |^2 a \right] + \dots = \sum_n \frac{|\xi|^n}{(2n)!} a - a + \frac{|\xi|^{2n}}{(2n+1)!} S^\dagger$$

$$-\xi |\xi|^2 a^\dagger$$

$$= a \operatorname{ch} |\xi| - a^\dagger \frac{\operatorname{sh} |\xi|}{|\xi|} = S^\dagger a S$$

$$S^\dagger a^\dagger S = (S^\dagger a S)^\dagger = a^\dagger \operatorname{ch} |\xi| - \xi^* a \operatorname{sh} |\xi|$$

$$\langle a_\varphi \rangle = \frac{\langle a \rangle e^{-i\varphi} + \langle a^\dagger \rangle e^{i\varphi}}{2} = \frac{\langle 0 | S^\dagger a S | 0 \rangle e^{-i\varphi}}{2} + \frac{e^{i\varphi} \langle 0 | S^\dagger a^\dagger S | 0 \rangle}{2}$$

$$= \langle 0 | a \cosh|\xi| - a^\dagger \frac{\sum sh|\xi|}{|\xi|} | 0 \rangle e^{-i\varphi} + \dots$$

$$\langle a_\varphi \rangle = 0 \Rightarrow \langle a_\varphi \rangle^2 = 0$$

$$\Delta a_\varphi^2 = \langle a_\varphi^2 \rangle = \frac{1}{2} \langle (a^\dagger e^{i\varphi} + a e^{-i\varphi})^2 \rangle = \frac{1}{2} (\langle a^{+2} \rangle e^{2i\varphi} + \langle a^2 \rangle e^{-2i\varphi} + \langle a^\dagger a \rangle + \langle a a^\dagger \rangle)$$

$$\langle a^2 \rangle = \langle 0 | S^\dagger a^2 S | 0 \rangle = \langle 0 | S^\dagger a S S^\dagger a S | 0 \rangle =$$

$$= \langle 0 | (a \cosh|\xi| - a^\dagger \frac{\sum sh|\xi|}{|\xi|})^2 | 0 \rangle =$$

$$= \langle 0 | a^2 \dots + a^{+2} \dots - a a^\dagger \frac{\sum ch|\xi| sh|\xi|}{|\xi|} - a^\dagger a \dots | 0 \rangle$$

$$= - \frac{\sum ch|\xi| sh|\xi|}{|\xi|} = \langle a^2 \rangle$$

$$\langle a^{+2} \rangle = \langle a^2 \rangle^* = - \frac{\sum^* ch|\xi| sh|\xi|}{|\xi|}$$

$$\langle a^\dagger a \rangle = \langle 0 | S^\dagger a^\dagger a S | 0 \rangle = \langle 0 | S^\dagger a^\dagger S S^\dagger a S | 0 \rangle =$$

↑ numero medio di fotoni del vuoto squeezed

$$= \langle 0 | (a^\dagger ch - \sum^* a^\dagger sh) (a ch - \sum a^\dagger sh) | 0 \rangle$$

$$= \langle 0 | a^{+2} \dots + a^2 \dots + a^\dagger a \dots + a a^\dagger \frac{|\xi|^2 sh^2 |\xi|}{|\xi|^2} | 0 \rangle$$

$$= \text{sh}^2 |\xi| = \langle a^\dagger a \rangle$$

$$\Delta a_\varphi^2 = \langle a_\varphi^2 \rangle = \frac{1}{2} \left( -\text{ch} |\xi| / \text{sh} |\xi| \frac{\sum^* e^{2i\varphi} + \sum e^{-2i\varphi}}{|\xi|} + 2 \frac{\text{sh}^2 |\xi|}{|\xi|} + 1 \right)$$

scelgo  $\varphi$  oppure  $\sum$  in modo t.c.

$$\frac{\sum^* e^{2i\varphi}}{|\xi|} = +1$$

$$\text{ch} x = \frac{e^x + e^{-x}}{2}$$

$$\Delta a_\varphi^2 = \frac{1}{2} \left( -2 \text{ch} |\xi| / \text{sh} |\xi| + 2 \text{sh}^2 |\xi| + 1 \right) = \text{sh} x = \frac{e^x - e^{-x}}{2}$$

$$= \frac{1}{2} \left( -2 \frac{e^x + e^{-x}}{2} \frac{e^x - e^{-x}}{2} + 2 \left( \frac{e^x - e^{-x}}{2} \right)^2 + 1 \right) =$$

$$= \frac{1}{2} e^{2|\xi|} = \Delta a_\varphi = \Delta p$$

stesso conto con  $\varphi + \frac{\pi}{2}$

$$\frac{\sum^* e^{2i\varphi}}{|\xi|} = +1 \Rightarrow \frac{\sum^* e^{2i(\varphi + \frac{\pi}{2})}}{|\xi|} = -1$$

$$\Delta a_{\varphi + \frac{\pi}{2}} = \Delta q^2 = e^{-2|\xi|}$$

$$\Delta a_\varphi = \Delta p = \frac{e^{|\xi|}}{\sqrt{2}}$$

$$\Delta q = \frac{e^{-|\xi|}}{\sqrt{2}}$$

quadratura  $e^-$ : il campo elettrico  $q = a_{\varphi=0}$  campo el

$$\vec{E} = -\frac{\partial \vec{A}}{\partial t}$$

evoluzione temporale della quadratura per evoluz

libera  $\rightarrow H = \hbar \omega (a^\dagger a + \frac{1}{2})$

$$a_\varphi(t) = e^{\frac{i}{\hbar} H t} a_\varphi e^{-\frac{i}{\hbar} H t}$$

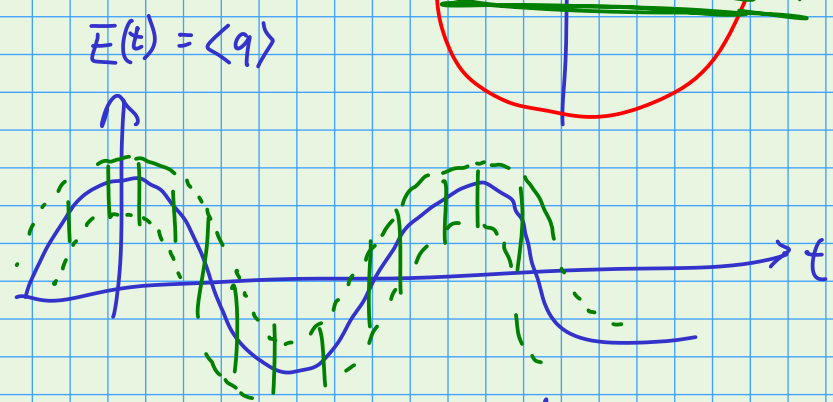
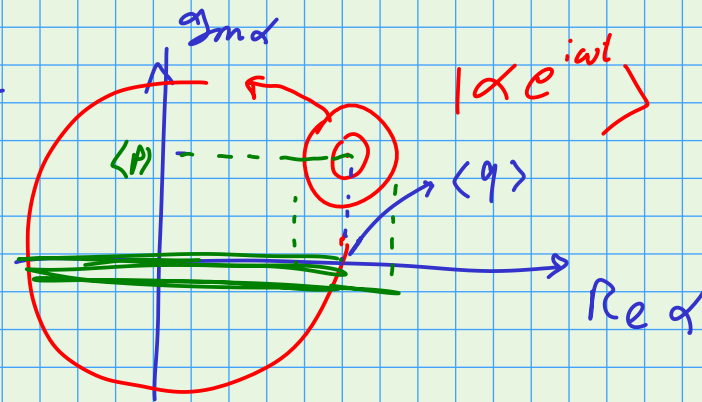


$$= e^{i\omega t a^\dagger a} a e^{-i\omega t a^\dagger a} = \frac{1}{\sqrt{2}} \left( e^{i\omega t a^\dagger a} a e^{-i\omega t a^\dagger a} e^{-i\varphi + hc} + \text{h.c.} \right)$$

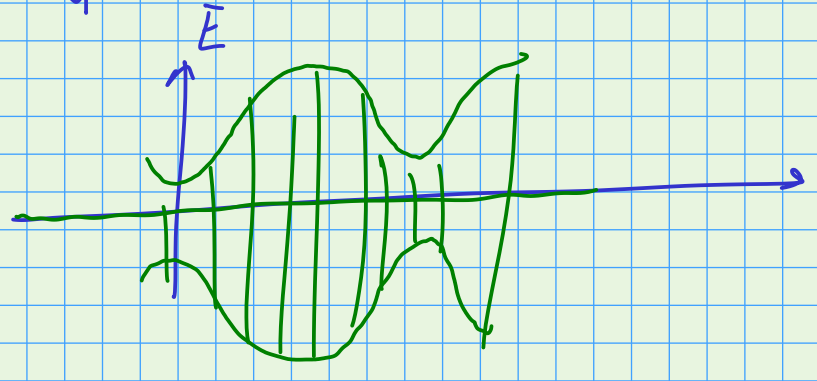
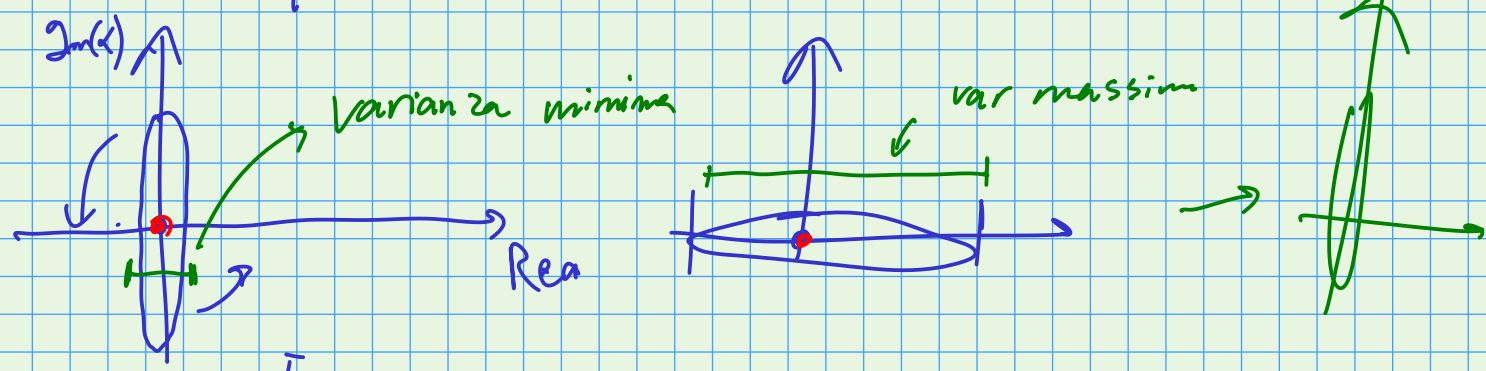
$$= \frac{1}{\sqrt{2}} \left( a^+ e^{i(\varphi + \omega t)} + a e^{-i(\varphi + \omega t)} \right) a e^{-i\omega t}$$

la quadratura ruota con periodo  $\omega$

Stato coerente



Vuoto squeezed





# coerente squeezed

